# Wave splitting for time-dependent scattered field separation 

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(1) Wave splitting

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(2) Wave splitting in the two-space dimensional case
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## Wave splitting

## Wave splitting

## Motivation

Time history of wave fields at one location: incident wave impinges on a sound-soft inclusion


## At location 1





## Wave splitting

Multiple scattering problem: $u=u_{1}+u_{2}, \quad$ in $\Omega:=\mathbb{R}^{d} \backslash\left(S_{1} \cup S_{2}\right)$

## $\Omega$


$u$ satifies:

$$
\frac{\partial^{2} u}{\partial t^{2}}-c^{2} \Delta u=0 \quad \text { in } \Omega, t>0
$$

Question: Given the measured total field $u$, can we recover $u_{1}$ and $u_{2}$ without knowing in advance either of them ?

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Outside $S_{1}$ and $S_{2}$, u satisfies:

$$
\frac{\partial^{2} u}{\partial t^{2}}-c^{2} \Delta u=0 \quad \text { in } \Omega, t>0
$$

$c>0$ constant.
At $t=0$, no signal in $\Omega$, then uniqueness of splitting ${ }^{1}$

$$
u=u_{1}+u_{2} \quad \text { in } \Omega, t>0
$$

and $u_{k}$ outgoing (3D):

$$
u_{k}\left(t, r_{k}, \theta_{k}, \varphi_{k}\right)=\frac{1}{r_{k}} \sum_{i \geq 0} \frac{f_{k, i}\left(r_{k}-c t, \theta_{k}, \varphi_{k}\right)}{\left(r_{k}\right)^{i}}
$$

$\left(r_{k}, \theta_{k}, \varphi_{k}\right)$ spherical coordinates centered at $C_{k}$.

[^0]Since

$$
u_{k}\left(t, r_{k}, \theta_{k}, \varphi_{k}\right)=\frac{1}{r_{k}} \sum_{i \geq 0} \frac{f_{k, i}\left(r_{k}-c t, \theta_{k}, \varphi_{k}\right)}{\left(r_{k}\right)^{i}}
$$

$\left(r_{k}, \theta_{k}, \varphi_{k}\right)$ spherical coordinates centered at $C_{k}$,
$m^{\mathrm{th}}$-order absorbing boundary condition ${ }^{2}$ on any $\Gamma$ in $\Omega$

$$
B_{k}\left[u_{k}\right]=O\left(\frac{1}{r_{k}^{2 m+1}}\right), \quad k=1,2
$$

[^1]
## Wave splitting

Neglecting the error term:

$$
B_{j}\left[u_{k}\right]=B_{j}\left[u_{k}+u_{j}\right]=B_{j}[u], \quad j=1,2, \quad k \neq j
$$

Recover $u_{1}$ and $u_{2}$ by solving:

$$
\begin{cases}B_{2}\left[u_{1}\right] & =B_{2}[u]  \tag{1}\\ B_{1}\left[u_{2}\right] & =B_{1}[u]\end{cases}
$$

where $u$ is known (measurements on $\Gamma$ )

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B_{j}\left[u_{k}\right]=B_{j}\left[u_{k}+u_{j}\right]=B_{j}[u], \quad j=1,2, \quad k \neq j
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Recover $u_{1}$ and $u_{2}$ by solving:

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\left\{\begin{array}{lll}
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B_{1}\left[u_{2}\right] & =B_{1}[u]
\end{array}\right.
$$

where $u$ is known (measurements on $\Gamma$ )
Difficulty: integration of partial differential equation (1)-(2) on the submanifold $\Gamma$

- Find adequate initial and boundary conditions
- Change of coordinates from $\left(r_{k}, \theta_{k}, \varphi_{k}\right)$ to $\left(r_{j}, \theta_{j}, \varphi_{j}\right)$
- Remove normal/radial derivatives (equation on $\Gamma$ involving only $\left.\left(t, \theta_{j}, \varphi_{j}\right)\right)$


## Wave splitting in the two-space dimensional case

 Change of coordinatesIn 2D, Bayliss-Turkel first order absorbing boundary condition

$$
B_{j}[u]=\frac{1}{c} \frac{\partial u}{\partial t}+\frac{\partial u}{\partial r_{j}}+\frac{u}{2 r_{j}}
$$

For simplicity, let $\Gamma:=\Gamma_{1} \cup \Gamma_{2}$ with

$$
\Gamma_{k}=\text { semi-circle centered at } C_{k}
$$



## Wave splitting in the two-space dimensional case

E.g. to recover $u_{1}$ on $\Gamma_{1}$ (semi-circle centered at $C_{1}$ )

$$
\begin{aligned}
B_{2}\left[u_{1}\right] & =B_{2}[u] \\
\frac{1}{c} \frac{\partial u_{1}}{\partial t}+\frac{\partial u_{1}}{\partial r_{2}}+\frac{u_{1}}{2 r_{2}} & =\frac{1}{c} \frac{\partial u}{\partial t}+\frac{\partial u}{\partial r_{2}}+\frac{u}{2 r_{2}}
\end{aligned}
$$

How to solve this PDE for $u_{1}$ ?

- need initial and boundary conditions
- remove the radial derivative! we solve on $\Gamma_{1}$
- derivatives in $\left(r_{2}, \theta_{2}\right)$, when domain in $\left(r_{1}, \theta_{1}\right)$
$\Longrightarrow$ rewrite the PDE using only $\frac{\partial}{\partial t}, \frac{\partial}{\partial \theta_{1}}$ and $0^{\text {th }}$-order term


## Wave splitting in the two-space dimensional case

 Change of coordinatesFinally: PDE to recover $f_{1}=\sqrt{r_{1}} u_{1}$ on $\Gamma_{1}, t>0$

$$
\left(\alpha_{1}\left(\theta_{1}\right) \frac{\partial}{\partial t}+\beta_{1}\left(\theta_{1}\right) \frac{\partial}{\partial \theta_{1}}+\gamma_{1}\left(\theta_{1}\right)\right) f_{1}=\left(\frac{1}{c} \frac{\partial}{\partial t}+\frac{\partial}{\partial r_{2}}+\frac{1}{2 r_{2}}\right) u
$$

with

$$
\begin{aligned}
\alpha_{1}\left(\theta_{1}\right) & =\frac{\sqrt{r_{1}^{2}+\ell^{2}-2 r_{1} \ell \cos \left(\theta_{1}\right)}-r_{1}+\ell \cos \left(\theta_{1}\right)}{c \sqrt{r_{1}} \sqrt{r_{1}^{2}+\ell^{2}-2 r_{1} \ell \cos \left(\theta_{1}\right)}}, \\
\beta_{1}\left(\theta_{1}\right) & =\frac{\ell \sin \left(\theta_{1}\right)}{r_{1} \sqrt{r_{1}} \sqrt{r_{1}^{2}+\ell^{2}-2 r_{1} \ell \cos \left(\theta_{1}\right)}}, \\
\gamma_{1}\left(\theta_{1}\right) & =\frac{\ell \cos \left(\theta_{1}\right)}{2 r_{1} \sqrt{r_{1}} \sqrt{r_{1}^{2}+\ell^{2}-2 r_{1} \ell \cos \left(\theta_{1}\right)}},
\end{aligned}
$$

and $\ell$ the signed distance between $C_{1}$ and $C_{2}$.

We want to recover $f_{1}=\sqrt{r_{1}} u_{1}$ which satisfies on $\Gamma, t>0$

$$
\left(\alpha_{1}\left(\theta_{1}\right) \frac{\partial}{\partial t}+\beta_{1}\left(\theta_{1}\right) \frac{\partial}{\partial \theta_{1}}+\gamma_{1}\left(\theta_{1}\right)\right) f_{1}=\left(\frac{1}{c} \frac{\partial}{\partial t}+\frac{\partial}{\partial r_{2}}+\frac{1}{2 r_{2}}\right) u
$$

## Initial condition?

At $t=0$, no signal in $\Omega$ : all sources in $S_{1} \cup S_{2}$
$\Longrightarrow f_{1}$ and $f_{2}$ vanish in $\Omega$, thus on $\Gamma_{1} \cup \Gamma_{2}$
the initial condition is:

$$
f_{1}=0, \quad \text { on } \Gamma_{1} \text {, at } t=0 .
$$

## Wave splitting in the two-space dimensional case

Hyperbolic PDE

$$
\left(\alpha_{1}\left(\theta_{1}\right) \frac{\partial}{\partial t}+\beta_{1}\left(\theta_{1}\right) \frac{\partial}{\partial \theta_{1}}+\gamma_{1}\left(\theta_{1}\right)\right) f_{1}=\left(\frac{1}{c} \frac{\partial}{\partial t}+\frac{\partial}{\partial r_{2}}+\frac{1}{2 r_{2}}\right) u
$$

trivial at $\theta_{1}=0$ or $\pi$ modulo $2 \pi$, since $\alpha_{1}\left(\theta_{1}\right)=0, \beta_{1}\left(\theta_{1}\right)=0$
$\Longrightarrow$ Dirichlet boundary condition: $f_{1}=\frac{B_{2}[u]}{\gamma_{1}(0)} \quad$ at $\theta_{1}=0$


Reconstruction of $f_{1}\left(\right.$ resp. $\left.f_{2}\right)$ on $\Gamma_{1}\left(\right.$ resp. $\left.\Gamma_{2}\right)$

by subtraction, $f_{2}\left(\right.$ resp. $\left.f_{1}\right)$ can be reconstructed on $\Gamma_{1}\left(\right.$ resp. $\left.\Gamma_{2}\right)$

## Numerical results in the 2D－case

## At location 2

Time history of wave fields at one location：two purely radial wave fields generated by point－sources





## Numerical results in the 2D－case

Time history of wave fields at one location：incident wave impinges on a sound－soft inclusion

## At location 1




## Conclusion

New partial differential equation

- on a submanifold 「
- local in space and time
- in the time-dependent domain

Method extendable to:

- 2 or more scatterers
- vector-valued wave equations from electromagnetics and elasticity
- improved accuracy with higher order absorbing boundary condition (more terms in the progressive wave expansion)
M.J. Grote, M. Kray, F. Nataf and F. Assous.

Wave splitting for time-dependent scattered field separation. C. R. Acad. Sci. Paris, Serie I (2015)


[^0]:    ${ }^{1}$ M. J. Grote and C. Kirsch. Nonreflecting boundary condition for time-dependent multiple scattering. J. Comput. Phys., 221(1):41-67, 2007.

[^1]:    ${ }^{2}$ A. Bayliss and E. Turkel. Radiation boundary conditions for wave-like equations. Comm. Pure Appl. Math., 33(6):707-725, 1980.

